

Problem 1:

a)

$$\nabla \cdot \vec{v} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = \frac{1}{1+t} + \frac{1}{1+2t} + 0 \neq 0$$

$$\boxed{\nabla \cdot \vec{v} = \frac{1}{1+t} + \frac{1}{1+2t}}$$

Why is $\nabla \cdot \vec{v} \neq 0$? Because the fluid is compressible and non-steady, and the continuity equation is:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0$$

$$\nabla \cdot \vec{v} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ u & v & w \end{vmatrix} = \begin{pmatrix} \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \\ \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \\ \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \end{pmatrix} = \vec{0}$$

$$\boxed{\nabla \wedge \vec{v} = \vec{0}}$$

If the flow is irrotational, then there is a potential function.

$$\vec{v} = \nabla \phi \Rightarrow u = \frac{\partial \phi}{\partial x} = \frac{x}{1+t} \Rightarrow \phi = \frac{1}{1+t} \cdot \frac{x^2}{2} + f(y, z, t)$$

$$v = \frac{\partial \phi}{\partial y} = \frac{\partial f}{\partial y} = \frac{y}{1+2t} \Rightarrow f = \frac{1}{2} \frac{y^2}{1+2t} + g(z, t)$$

$$w = \frac{\partial \phi}{\partial z} = \frac{\partial g}{\partial z} = 0 \Rightarrow g = C(t)$$

$$\boxed{\phi = \frac{1}{1+t} \cdot \frac{x^2}{2} + \frac{1}{1+2t} \frac{y^2}{2} + C(t)}$$

b) Streamlines: $\frac{dx}{u} = \frac{dy}{v} = \frac{dz}{w}$

$$\frac{dx}{u} = \frac{dx}{x/(1+t)} = \frac{dy}{y/(1+2t)} = \frac{dy}{v} \Rightarrow \frac{dy}{y} = \frac{1+t}{1+2t} \cdot \frac{dx}{x} \Rightarrow$$

$$\Rightarrow \ln(y) = \ln(x)^{1+t/1+2t} + \ln(C_1) \Rightarrow$$

$$\Rightarrow \boxed{y = C_1(t) \cdot x^{1+t/1+2t}}$$

$$\frac{dz}{w} = \frac{dy}{v} \Rightarrow dz = 0 \Rightarrow \boxed{z = \text{const}}$$

Trajectories: $\vec{x}(t=t_0) = \vec{x}_0 = (x_0, y_0)$

$$\frac{dx}{dt} = u = \frac{x}{1+t} \Rightarrow \frac{dx}{x} = \frac{dt}{1+t} \Rightarrow \ln(x) = \ln(1+t) + \ln(C_1) \Rightarrow$$

$$\Rightarrow x = C_1 \cdot (1+t)$$

$$\frac{dy}{dt} = v = \frac{y}{1+2t} \Rightarrow \frac{dy}{y} = \frac{dt}{1+2t} \Rightarrow \ln(y) = \frac{1}{2} \ln(1+2t) + \ln(C_2) \Rightarrow$$

$$\Rightarrow y = C_2 \cdot (1+2t)^{1/2}$$

$$\frac{dz}{dt} = w = 0 \Rightarrow z = C_3$$

$$t = t_0 \Rightarrow \begin{cases} x_0 = C_1 \\ y_0 = C_2 \\ z_0 = C_3 \end{cases} \Rightarrow \left\{ \begin{array}{l} x = x_0(1+t) \\ y = y_0(1+2t)^{1/2} \\ z = z_0 \end{array} \right.$$

$$\boxed{\begin{array}{l} x = x_0(1+t) \\ y = y_0(1+2t)^{1/2} \\ z = z_0 \end{array}}$$

Paths: Eliminate t

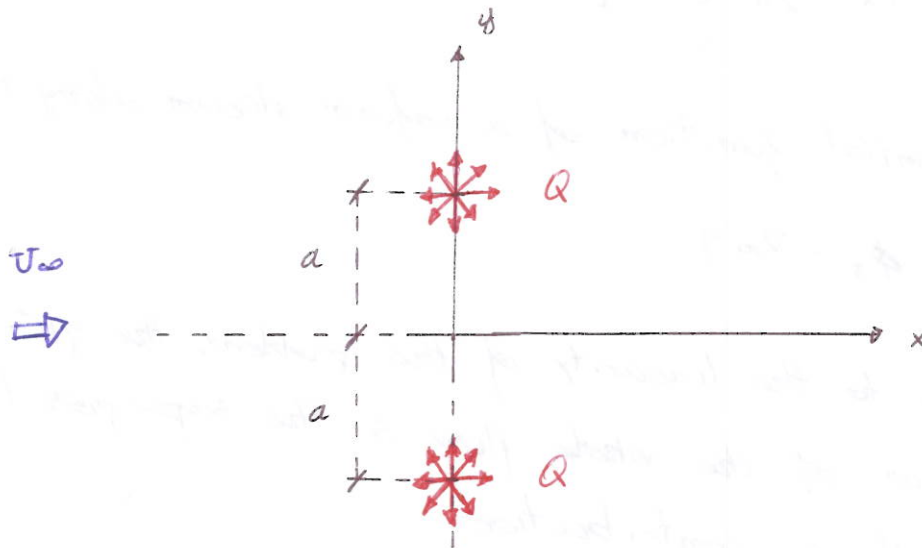
$$t = \frac{x}{x_0} - 1 \Rightarrow$$

$$\boxed{\begin{array}{l} y = y_0 \left(\frac{x}{x_0} - 1 \right)^{1/2} \\ z = z_0 \end{array}}$$

c) streaklines: fluid particles that in previous times ($t = \tau$) passed by the point $\vec{x}_p = (x_p, y_p)$

$$x_p = C_1 \cdot (1 + \tau) \Rightarrow C_1 = \frac{x_p}{1 + \tau} \Rightarrow \boxed{x = x_p \cdot \frac{1 + t}{1 + \tau}}$$

$$y_p = C_2 \cdot (1 + 2\tau)^{1/2} \Rightarrow C_2 = \frac{y_p}{(1 + 2\tau)^{1/2}} \Rightarrow \boxed{y = y_p \cdot \left(\frac{1 + 2t}{1 + 2\tau} \right)^{1/2}}$$

Problem 2:

1)

* In order to calculate the potential function for the upper source ϕ_1 , we follow these steps:

- Potential function for a source at the origin in polar coordinates:

$$\phi_0 = \frac{Q}{2\pi} \ln\left(\frac{r}{a}\right)$$

- In cartesian coordinates:

$$\phi_0 = \frac{Q}{2\pi} \ln\left[\frac{(x^2 + y^2)^{1/2}}{a}\right] = \frac{Q}{2\pi} \frac{1}{2} \ln\left[\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a}\right)^2\right]$$

- translate the source to $(x, y) = (0, a)$:

$$\phi_1 = \frac{Q}{2\pi} \frac{1}{2} \ln\left[\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a} - 1\right)^2\right]$$

* Repeat with the source located at $(x, y) = (0, -a)$:

$$\phi_2 = \frac{Q}{2\pi} \frac{1}{2} \ln \left[\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} - 1 \right)^2 \right]$$

* Potential function of a uniform stream along the x-axis:

$$\phi_3 = v_\infty x$$

* Due to the linearity of the problem, the potential function of the whole flow is the superposition of the three contributions:

$$\phi(x, y) = v_\infty x + \frac{Q}{2\pi} \cdot \frac{1}{2} \left\{ \ln \left[\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} - 1 \right)^2 \right] + \ln \left[\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} + 1 \right)^2 \right] \right\}$$

$$2) \quad * \quad u = \frac{\partial \phi}{\partial x} = \frac{\partial \psi}{\partial y} = v_\infty + \frac{Q}{2\pi} \frac{1}{2} \left\{ \frac{1}{a} \frac{2x/a}{\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} - 1 \right)^2} + \frac{1}{a} \frac{2x/a}{\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} + 1 \right)^2} \right\} \Rightarrow$$

$$\Rightarrow u(x, y) = v_\infty + \frac{Q}{2\pi} \left\{ \frac{x}{x^2 + (y-a)^2} + \frac{x}{x^2 + (y+a)^2} \right\}$$

$$* \quad v = \frac{\partial \phi}{\partial y} = -\frac{\partial \psi}{\partial x} = \frac{Q}{2\pi} \frac{1}{2} \left\{ \frac{1}{a} \frac{2(\frac{y}{a} - 1)}{\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} - 1 \right)^2} + \frac{1}{a} \frac{2(\frac{y}{a} + 1)}{\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} + 1 \right)^2} \right\} \Rightarrow$$

$$\Rightarrow v(x, y) = \frac{Q}{2\pi} \left\{ \frac{y-a}{x^2 + (y-a)^2} + \frac{y+a}{x^2 + (y+a)^2} \right\}$$

3) The compatibility conditions relate the streamfunction and the potential function:

$$\begin{cases} u = \frac{\partial \psi}{\partial y} = \frac{\partial \phi}{\partial x} \\ v = -\frac{\partial \psi}{\partial x} = \frac{\partial \phi}{\partial y} \end{cases}$$

Applying the first:

$$* u = \frac{\partial \phi}{\partial x} = \frac{\partial \psi}{\partial y} \Rightarrow \psi = \int u dy = v_{\infty} y + \frac{Q}{2\pi} \left\{ a \tan\left(\frac{y-a}{x}\right) + a \tan\left(\frac{y+a}{x}\right) \right\} + F(x)$$

We need to calculate the free function $F(x)$. We can do it by plugging $\psi(x,y)$ in the second compatibility condition:

$$* v = \frac{Q}{2\pi} \left\{ \frac{y-a}{x^2+(y-a)^2} + \frac{y+a}{x^2+(y+a)^2} \right\} = -\frac{\partial \psi}{\partial x} = \frac{Q}{2\pi} \left\{ \frac{y-a}{x^2+(y-a)^2} + \frac{y+a}{x^2+(y+a)^2} \right\} - F'(x) \Rightarrow$$

$$\Rightarrow F'(x) = 0 \Rightarrow F(x) = \text{cte.}$$

As expected, the streamfunction (and the potential function) are defined up to a free constant. We can choose any value that we want. I choose $F(y) = 0$ because it makes things simpler:

$$* \boxed{\psi(x,y) = v_{\infty} y + \frac{Q}{2\pi} \left\{ a \tan\left(\frac{y-a}{x}\right) + a \tan\left(\frac{y+a}{x}\right) \right\}}$$

4)

$$* \frac{u}{Q/2\pi a} = \frac{U_\infty}{Q/2\pi a} + \frac{Q/2\pi}{Q/2\pi a} \left\{ \frac{1}{a} \frac{r \cos \theta / a}{\left(\frac{r \cos \theta}{a}\right)^2 + \left(\frac{r \sin \theta}{a} - 1\right)^2} + \frac{1}{a} \frac{r \cos \theta / a}{\left(\frac{r \cos \theta}{a}\right)^2 + \left(\frac{r \sin \theta}{a} + 1\right)^2} \right\} =$$

$$= K + \frac{R \cos \theta}{R^2 \cos^2 \theta + (R \sin \theta - 1)^2} + \frac{R \cos \theta}{R^2 \cos^2 \theta + (R \sin \theta + 1)^2} =$$

$$= K + R \cos \theta \left\{ \frac{1}{R^2 + 2R \sin \theta + 1} + \frac{1}{R^2 - 2R \sin \theta + 1} \right\} \Rightarrow$$

$$\Rightarrow \frac{u(R, \theta)}{Q/2\pi a} = K + \frac{2R \cos \theta (R^2 + 1)}{(R^2 + 1)^2 - 4R^2 \sin^2 \theta}$$

$$* \frac{v}{Q/2\pi a} = \frac{Q/2\pi}{Q/2\pi a} \left\{ \frac{1}{a} \frac{\frac{r \sin \theta}{a} - 1}{\left(\frac{r \cos \theta}{a}\right)^2 + \left(\frac{r \sin \theta}{a} - 1\right)^2} + \frac{1}{a} \frac{\frac{r \sin \theta}{a} + 1}{\left(\frac{r \cos \theta}{a}\right)^2 + \left(\frac{r \sin \theta}{a} + 1\right)^2} \right\} =$$

$$= \frac{R \sin \theta - 1}{R^2 \cos^2 \theta + (R \sin \theta - 1)^2} + \frac{R \sin \theta + 1}{R^2 \cos^2 \theta + (R \sin \theta + 1)^2} =$$

$$= \frac{R \sin \theta - 1}{R^2 - 2R \sin \theta + 1} + \frac{R \sin \theta + 1}{R^2 + 2R \sin \theta + 1} \Rightarrow$$

$$\Rightarrow \frac{v(R, \theta)}{Q/2\pi a} = \frac{2R \sin \theta (R^2 - 1)}{(R^2 + 1)^2 - 4R^2 \sin^2 \theta}$$

5)

$$* v(R, \theta) = 0 \Rightarrow 2R \sin \theta (R^2 - 1) = 0 \Rightarrow$$

$$R = 0$$

$$R^2 - 1 = 0 \Rightarrow$$

$$R = 1$$

$$\sin \theta = 0 \Rightarrow$$

$$\theta = 0$$

$$\theta = \pi$$

6)

* $R=0$:

$$u(R=0, \theta) = 0 \Rightarrow K=0 \Rightarrow \text{Only valid if } u_0=0$$

* $R=1$:

$$u(R=1, \theta) = 0 \Rightarrow K + \frac{2 \cos \theta \cdot 2}{4 - 4 \sin^2 \theta} = K + \frac{\cos \theta}{1 - \sin^2 \theta} = K + \frac{\cos \theta}{\cos^2 \theta} =$$

$$= K + \frac{1}{\cos \theta} = 0 \Rightarrow \cos \theta = -\frac{1}{K}$$

$$\Rightarrow K > 1 \Rightarrow \cos \theta = -\frac{1}{K} \Rightarrow \theta = \arccos\left(-\frac{1}{K}\right) + (2n+1)\pi$$

$$\Rightarrow K = 1 \Rightarrow \cos \theta = -1 \Rightarrow \theta = \arccos(-1) = (2n+2)\pi$$

$$\Rightarrow K < 1 \Rightarrow \cos \theta = -\frac{1}{K} < -1 \Rightarrow \text{No solution}$$

* $\theta=0$:

$$u(R, \theta=0) = 0 \Rightarrow K + \frac{2R(R^2+1)}{(R^2+1)^2} = K + \frac{2R}{R^2+1} = 0 \Rightarrow$$

$$\Rightarrow KR^2 + 2R + K = 0 \Rightarrow R = \frac{-2 \pm \sqrt{4 - 4K^2}}{2K} = \frac{-1 \pm \sqrt{\frac{1}{K^2} - 1}}{K}$$

$$\text{Need } R \geq 0 \Rightarrow R = -\frac{1}{K} + \sqrt{\frac{1}{K^2} - 1}$$

$$\Rightarrow K > 1 \Rightarrow \frac{1}{K^2} - 1 < 0 \Rightarrow \text{No physical solution.}$$

$$\Rightarrow K = 1 \Rightarrow R = -1 \Rightarrow \text{No physical solution.}$$

$$\Rightarrow K < 1 \Rightarrow R = -\frac{1}{K} + \sqrt{\frac{1}{K^2} - 1} < 0 \Rightarrow \text{No physical solution.}$$

* $\theta = \pi$:

$$\bullet u(R, \theta = \pi) = 0 \Rightarrow K + \frac{2R(-1)(R^2+1)}{(R^2+1)^2 - 4R^2 \cdot 0} = K - \frac{2R}{R^2+1} = 0 \Rightarrow$$

$$\Rightarrow KR^2 - 2R + K = 0 \Rightarrow R = \frac{2 \pm \sqrt{4 - 4K^2}}{2K} = \frac{1}{K} \pm \sqrt{\frac{1}{K^2} - 1}$$

$$\Rightarrow K > 1 \Rightarrow \frac{1}{K^2} - 1 < 0 \Rightarrow \text{No physical solution.}$$

$$\Rightarrow K = 1 \Rightarrow R = 1 \pm \sqrt{0} = 1$$

$$\Rightarrow K < 1 \Rightarrow R = \frac{1}{K} \pm \sqrt{\frac{1}{K^2} - 1} > 0$$

* Summary:

$\Rightarrow \underline{K < 1}$: two simple stagnation points.

$$\bullet (R, \theta)_1 = \left(\frac{1}{K} + \sqrt{\frac{1}{K^2} - 1}, \pi \right)$$

$$\bullet (R, \theta)_2 = \left(\frac{1}{K} - \sqrt{\frac{1}{K^2} - 1}, \pi \right)$$

$\Rightarrow \underline{K = 1}$: one double stagnation point.

$$\bullet (R, \theta)_1 = (1, \pi)$$

$$\bullet (R, \theta)_2 = (1, \pi)$$

$\Rightarrow \underline{K > 1}$:

$$\bullet (R, \theta)_1 = \left[1, \arccos\left(-\frac{1}{K}\right) \right]$$

$$\bullet (R, \theta)_2 = \left[1, -\arccos\left(-\frac{1}{K}\right) + 2\pi \right]$$

7)

$$\Psi(x, y) = K \frac{y}{a} + \left\{ a \tan^{-1} \left(\frac{y/a - 1}{x/a} \right) + a \tan^{-1} \left(\frac{y/a + 1}{x/a} \right) \right\} \frac{Q/2\pi}{}$$

* $K \ll 1$:

• Stagnation Point 1: $(R, \theta)_1 = (R, \pi) \Rightarrow (x, y)_1 = (-R \cdot a, 0)$
LO

* In every point of a streamline, the value of $\Psi(x, y)$ is constant. In the streamline that passes through the stagnation point 1, $\Psi_1(x, y)$, the value is:

$$\frac{\Psi_1(x_1, y_1)}{Q/2\pi} = a \tan^{-1} \left(\frac{-1}{-R_1} \right) + a \tan^{-1} \left(\frac{1}{-R_1} \right) = 0$$

* When $x \rightarrow \pm \infty$:

$$\lim_{x \rightarrow \pm \infty} \left[a \tan^{-1} \left(\frac{y/a - 1}{x/a} \right) \right] = \pm n \frac{\pi}{2}$$

$$\lim_{x \rightarrow \pm \infty} \left[\frac{\Psi(x, y)}{Q/2\pi} \right] = K \cdot \frac{y}{a} \pm n\pi = \frac{\Psi_1(x_1, y_1)}{Q/2\pi} = 0$$

* The asymptotes of the streamlines that pass through the Stagnation Point 1 are:

$$K \frac{y}{a} \pm n\pi = 0 \Rightarrow y = \begin{cases} -\pi a / K \\ 0 \\ +\pi a / K \end{cases}$$

- Stagnation Point 2: $(R_1, \theta)_2 = (R_2, \pi) \Rightarrow (x, y)_2 = (-R_2 \cdot a, 0)$
 $\angle 0$

* Streamline that passes through the stagnation Point 2 ($\Psi_2(x, y)$):

- $\frac{\Psi_2(x_2, y_2)}{Q/2\pi} = a \tan\left(\frac{-1}{-R_2}\right) + a \tan\left(\frac{1}{-R_2}\right) = 0$

* When $x \rightarrow \pm \infty$:

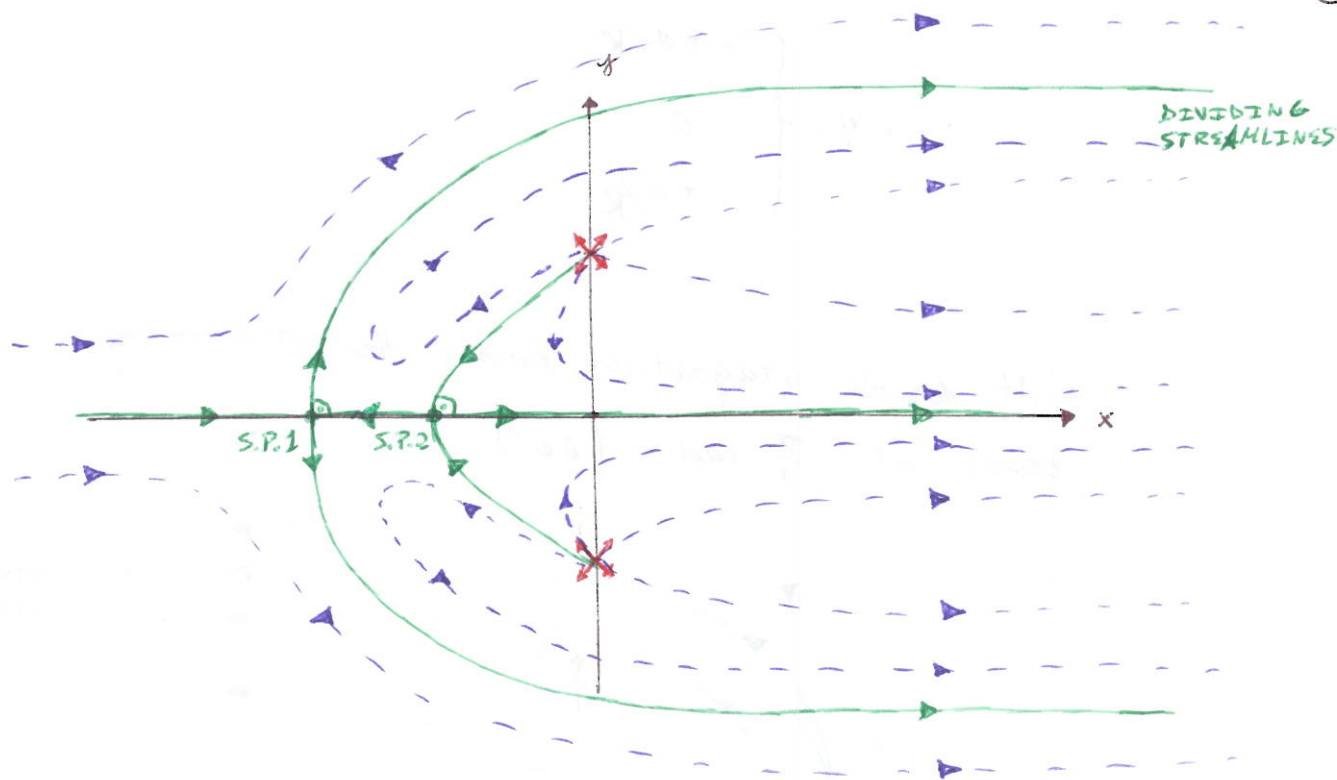
- $\lim_{x \rightarrow \pm \infty} \left[a \tan\left(\frac{\frac{y}{a} - 1}{x/a}\right) \right] = \pm n \cdot \frac{\pi}{2}$

- $\lim_{x \rightarrow \pm \infty} \left[\frac{\Psi_2(x, y)}{Q/2\pi} \right] = K \cdot \frac{y}{a} \pm n\pi = \frac{\Psi_2(x_2, y_2)}{Q/2\pi} = 0$

* The asymptotes of the streamlines that pass through the stagnation Point 2 are:

- $y = \begin{cases} -\pi a/K \\ 0 \\ +\pi a/K \end{cases}$

• The streamlines only cross at singular points (sources, sinks, doublets, etc) and stagnation points. At simple stagnation points, they cross at $\frac{\pi}{2}$ rad (90°).



* $K = 1$:

• Double Stagnation Point: $(R, \theta)_1 = (1, \pi) \Rightarrow (x, y)_1 = (-a, 0)$

* Streamline that passes through the stagnation point $(\psi_1(x, y))$:

$$\bullet \frac{\psi_1(x_1, y_1)}{Q/2\pi} = a \tan\left(\frac{1}{a}\right) + a \tan\left(\frac{1}{-a}\right) = 0$$

* When $x \rightarrow \pm\infty$:

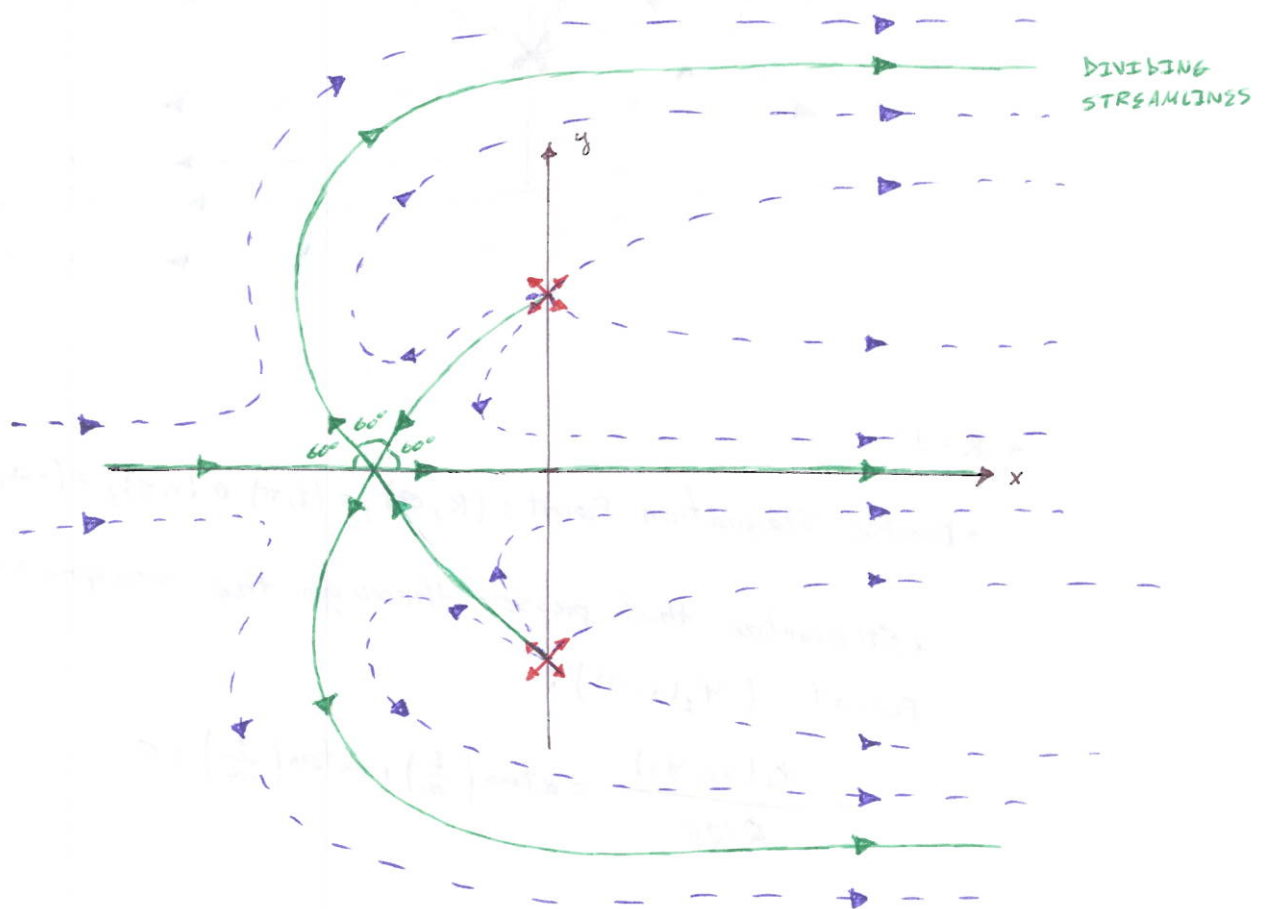
$$\bullet \lim_{x \rightarrow \pm\infty} \left[a \tan\left(\frac{\frac{y}{a} - 1}{x/a}\right) \right] = \pm n \frac{\pi}{2}$$

$$\bullet \lim_{x \rightarrow \pm\infty} \left[\frac{\psi_1(x, y)}{Q/2\pi} \right] = K \frac{y}{a} \pm n\pi = \frac{\psi_1(x_1, y_1)}{Q/2\pi} = 0$$

* The asymptotes of the streamlines that pass through the stagnation point are:

$$y = \begin{cases} -\pi a/K \\ 0 \\ +\pi a/K \end{cases}$$

- At double stagnation points, the streamlines cross at $\frac{\pi}{3}$ rad (60°).



* $K > 1$:

- Stagnation Point 1: $(R, \theta)_1 = [1, a \cos(\frac{1}{K})] \Rightarrow (x, y)_1 = (-\frac{a}{K}, a \sqrt{1 - \frac{1}{K^2}})$

* Streamline that passes through the stagnation Point 1 ($\psi_1(x, y)$):

$$\frac{\psi_1(x, y)}{Q/2\pi} = \sqrt{K^2 - 1} + \left[a \tan(K - \sqrt{K^2 - 1}) + a \ln(-K - \sqrt{K^2 - 1}) \right]$$

* When $x \rightarrow \pm \infty$:

$$\lim_{x \rightarrow \pm \infty} \left[\frac{\psi_1(x, y)}{Q/2\pi} \right] = K \frac{y}{a} \pm n\pi$$

* The asymptotes of the streamlines that pass through the stagnation Point 1 are:

$$y = \begin{cases} \psi_1 / a v_\infty \\ \frac{\psi_1}{a v_\infty} + \frac{\pi}{K} \end{cases}$$

$$\bullet \text{ Stagnation Point 2: } (R, \theta)_2 = \left[1, -a \cos\left(\frac{1}{K}\right) \right] \Rightarrow (x, y)_2 = \left(\frac{-a}{K}, -a \sqrt{1 - \frac{1}{K^2}} \right)$$

* Streamline that passes through the stagnation Point 2 ($\psi_2(x, y)$):

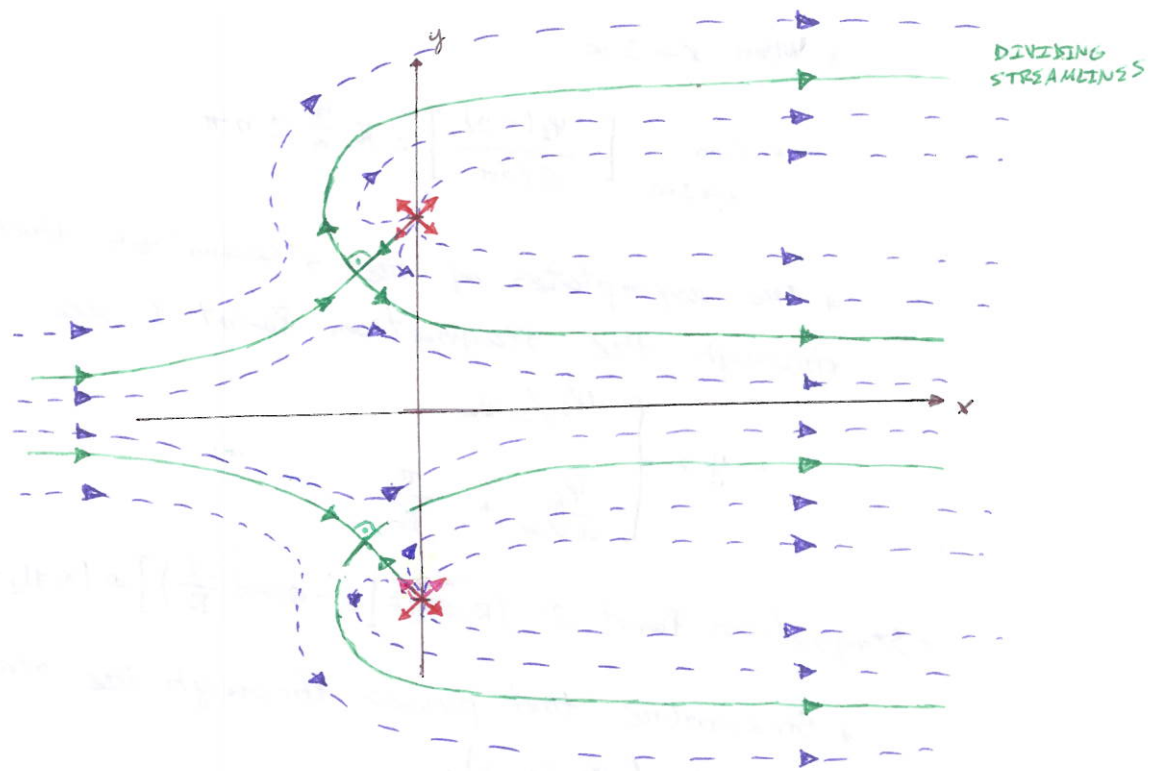
$$\bullet \frac{\psi_2(x_2, y_2)}{Q/2\pi} = -\sqrt{K^2 - 1} + \left\{ a \tan(K + \sqrt{K^2 - 1}) + a \tan(-K - \sqrt{K^2 - 1}) \right\}$$

* When $x \rightarrow \pm \infty$:

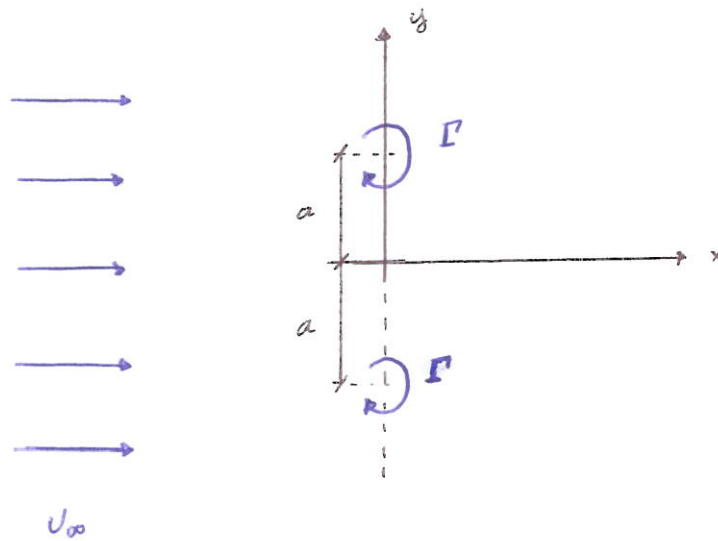
$$\lim_{x \rightarrow \pm \infty} \left[\frac{\psi_2(x, y)}{Q/2\pi} \right] = K \frac{y}{a} \pm n\pi$$

* The asymptotes of the streamlines that pass through the stagnation Point 2 are:

$$y = \begin{cases} \frac{\psi_2}{a v_\infty} \\ \frac{\psi_2}{a v_\infty} - \frac{\pi}{K} \end{cases}$$



Problem 1:



1)

⇒ stream function for a vortex at the origin in polar coordinates: $\psi_0 = \frac{\Gamma}{2\pi} \ln\left(\frac{r}{a}\right)$

⇒ In Cartesian coordinates: $\psi_0 = \frac{\Gamma}{2\pi} \ln\left[\frac{(x^2 + y^2)^{1/2}}{a}\right] = \frac{\Gamma}{2\pi} \cdot \frac{1}{2} \ln\left[\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a}\right)^2\right]$

⇒ If the vortex is located at $(x, y) = (0, a)$:

$$\psi_1 = \frac{\Gamma}{2\pi} \cdot \frac{1}{2} \cdot \ln\left[\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a} - 1\right)^2\right]$$

⇒ For a vortex located at $(x, y) = (0, -a)$:

$$\psi_2 = \frac{\Gamma}{2\pi} \cdot \frac{1}{2} \cdot \ln\left[\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a} + 1\right)^2\right]$$

⇒ stream function of a uniform stream along the x-axis:

$$\psi_3 = U_\infty y$$

⇒ the stream function for the whole flow is the superposition of the three of them:

$$\psi(x, y) = U_\infty y + \frac{\Gamma}{4\pi} \cdot \left\{ \ln\left[\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a} + 1\right)^2\right] + \ln\left[\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a} - 1\right)^2\right] \right\}$$

2)

$$u = \frac{\partial \psi}{\partial y} = v_{\infty} + \frac{\Gamma}{4\pi} \left\{ 2 \cdot \frac{y+a}{x^2+(y+a)^2} + 2 \cdot \frac{y-a}{x^2+(y-a)^2} \right\} \Rightarrow$$

$$\Rightarrow \boxed{u(x, y) = v_{\infty} + \frac{\Gamma}{2\pi} \left\{ \frac{y+a}{x^2+(y+a)^2} + \frac{y-a}{x^2+(y-a)^2} \right\}}$$

$$v = -\frac{\partial \psi}{\partial x} = \frac{-\Gamma}{4\pi} \cdot \left\{ \frac{2x}{x^2+(y+a)^2} + \frac{2x}{x^2+(y-a)^2} \right\} \Rightarrow$$

$$\Rightarrow \boxed{v(x, y) = \frac{-\Gamma}{2\pi} \left\{ \frac{x}{x^2+(y+a)^2} + \frac{x}{x^2+(y-a)^2} \right\}}$$

$$3) \quad u = \frac{\partial \psi}{\partial y} = \frac{\partial \phi}{\partial x} \Rightarrow \phi = \int u \, dx = \int \left\{ v_{\infty} + \frac{\Gamma}{2\pi} \left[\frac{y+a}{x^2+(y+a)^2} + \frac{y-a}{x^2+(y-a)^2} \right] \right\} dx =$$

$$= v_{\infty} x + \frac{\Gamma}{2\pi} \left\{ \arctan\left(\frac{x}{y+a}\right) + \arctan\left(\frac{x}{y-a}\right) \right\} + F(y)$$

$$v = -\frac{\partial \psi}{\partial x} = \frac{-\Gamma}{2\pi} \left\{ \frac{x}{x^2+(y+a)^2} + \frac{x}{x^2+(y-a)^2} \right\} = \frac{\partial \phi}{\partial y} =$$

$$= \frac{\Gamma}{2\pi} \cdot \left\{ \frac{-x}{x^2+(y+a)^2} - \frac{x}{x^2+(y-a)^2} \right\} + F'(y) \Rightarrow F'(y) = 0 \Rightarrow$$

$$\Rightarrow F(y) = \text{cte.}$$

$$\boxed{\phi(x, y) = v_{\infty} \cdot x + \frac{\Gamma}{2\pi} \left\{ \arctan\left(\frac{x}{y+a}\right) + \arctan\left(\frac{x}{y-a}\right) \right\}}$$

4)

$$\begin{aligned} \frac{U}{I/2\pi a} &= \frac{U_{00}}{I/2\pi a} + \left\{ \frac{\frac{y}{a} + 1}{\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a} + 1\right)^2} + \frac{\frac{y}{a} - 1}{\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a} - 1\right)^2} \right\} = \\ &= K + \frac{R \sin \theta + 1}{(R \cos \theta)^2 + (R \sin \theta + 1)^2} + \frac{R \sin \theta - 1}{(R \cos \theta)^2 + (R \sin \theta - 1)^2} = \\ &= K + \frac{R \sin \theta + 1}{R^2 + 2R \sin \theta + 1} + \frac{R \sin \theta - 1}{R^2 - 2R \sin \theta + 1} \Rightarrow \end{aligned}$$

$$\Rightarrow \frac{U}{I/2\pi a} = K + \frac{2R \sin \theta (R^2 - 1)}{(R^2 + 1)^2 - 4R^2 \sin^2 \theta}$$

$$\begin{aligned} \frac{V}{I/2\pi a} &= \frac{-\frac{x}{a}}{\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a} + 1\right)^2} - \frac{\frac{x}{a}}{\left(\frac{x}{a}\right)^2 + \left(\frac{y}{a} - 1\right)^2} = \\ &= \frac{-R \cos \theta}{(R \cos \theta)^2 + (R \sin \theta + 1)^2} - \frac{R \cos \theta}{(R \cos \theta)^2 + (R \sin \theta - 1)^2} = \\ &= -R \cos \theta \cdot \left\{ \frac{1}{R^2 + 2R \sin \theta + 1} + \frac{1}{R^2 - 2R \sin \theta + 1} \right\} \Rightarrow \end{aligned}$$

$$\Rightarrow \frac{V}{I/2\pi a} = \frac{-2R \cos \theta (R^2 + 1)}{(R^2 + 1)^2 - 4R^2 \sin^2 \theta}$$

$$5) \quad \frac{V}{I/2\pi a} = \frac{-2R \cos \theta (R^2 + 1)}{(R^2 + 1)^2 - 4R^2 \sin^2 \theta} = 0 \Rightarrow \begin{cases} R^2 = -1 \Rightarrow \text{NO!} \\ R = 0 \\ \cos \theta = 0 \Rightarrow \begin{cases} \theta = \pi/2 \\ \theta = 3\pi/2 \end{cases} \end{cases}$$

6)

$$\frac{u(R=0)}{I/2\pi a} = K = 0 \rightarrow \text{No solution unless } u_\infty = 0$$

$$\frac{u(\theta = \frac{\pi}{2})}{I/2\pi a} \stackrel{\sin(\frac{\pi}{2})=1}{=} K + \frac{2R(R^2-1)}{(R^2+1)^2-4R^2} = K + \frac{2R(R^2-1)}{R^4+2R^2+1-4R^2} =$$

$$= K + \frac{2R(R^2-1)}{(R^2-1)^2} = K + \frac{2R}{R^2-1} = 0 \Rightarrow$$

$$\Rightarrow KR^2 - K + 2R = 0 \Rightarrow R = \frac{-2 \pm \sqrt{4+4K^2}}{2K} \Rightarrow$$

$$\Rightarrow R = -\frac{1}{K} \pm \sqrt{1 + \frac{1}{K^2}}$$

$$R \text{ must be greater than } 0 \Rightarrow R = -\frac{1}{K} + \sqrt{1 + \frac{1}{K^2}}$$

$$\frac{u(\theta = \frac{3\pi}{2})}{I/2\pi a} \stackrel{\sin(\frac{3\pi}{2})=-1}{=} K + \frac{2R(-1)(R^2-1)}{(R^2+1)^2-4R^2} = K - \frac{2R}{R^2-1} = 0 \Rightarrow$$

$$\Rightarrow KR^2 - 2R - K = 0 \Rightarrow R = \frac{2 \pm \sqrt{4+4K^2}}{2K} \Rightarrow$$

$$\Rightarrow R = \frac{1}{K} \pm \sqrt{1 + \frac{1}{K^2}} \Rightarrow R = \frac{1}{K} + \sqrt{1 + \frac{1}{K^2}}$$

So the two stagnation points are:

$$(R, \theta)_1 = \left[-\frac{1}{K} + \sqrt{1 + \frac{1}{K^2}}, \frac{\pi}{2} \right]$$

$$(R, \theta)_2 = \left[\frac{1}{K} + \sqrt{1 + \frac{1}{K^2}}, \frac{3\pi}{2} \right]$$

In Cartesian coordinates:

$$(x_1, y_1) = \left(0, -\frac{1}{K} + \sqrt{1 + \frac{1}{K^2}} \right) \cdot a$$

$$(x_2, y_2) = \left(0, -\frac{1}{K} - \sqrt{1 + \frac{1}{K^2}} \right) \cdot a$$

7)

A streamline that passes through a stagnation point is a dividing streamline.

The non-dimensional stream function is:

$$\frac{\Psi(x, y)}{\Gamma/2\pi} = K \frac{y}{a} + \frac{1}{2} \left\{ \ln \left[\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} + 1 \right)^2 \right] + \ln \left[\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} - 1 \right)^2 \right] \right\}$$

The value of the stream function for the streamline that passes by the first stagnation point is:

$$\begin{aligned} \frac{\Psi_1(x_1, y_1)}{\Gamma/2\pi} &= K \frac{y_1}{a} + \frac{1}{2} \left\{ \ln \left[\left(\frac{x_1}{a} \right)^2 + \left(\frac{y_1}{a} + 1 \right)^2 \right] + \ln \left[\left(\frac{x_1}{a} \right)^2 + \left(\frac{y_1}{a} - 1 \right)^2 \right] \right\} = \\ &= K \frac{y_1}{a} + \frac{1}{2} \ln \left\{ \left(\frac{y_1}{a} + 1 \right)^2 \left(\frac{y_1}{a} - 1 \right)^2 \right\} = \\ &= K \frac{y_1}{a} + \frac{1}{2} \ln \left[\left(\frac{y_1}{a} \right)^2 - 1 \right]^2 = \\ &= -1 + \sqrt{K^2 + 1} + \frac{1}{2} \ln \left\{ \left[\left(-\frac{1}{K} + \sqrt{1 + \frac{1}{K^2}} \right)^2 - 1 \right]^2 \right\} = \\ &= -1 + \sqrt{K^2 + 1} + \ln \left\{ \frac{2}{K^2} - \frac{2}{K^2} \sqrt{K^2 + 1} \right\} \end{aligned}$$

and for the second:

$$\begin{aligned} \frac{\Psi_2(x_2, y_2)}{\Gamma/2\pi} &= K \frac{y_2}{a} + \ln \left[\left(\frac{y_2}{a} \right)^2 - 1 \right] = \\ &= -1 - \sqrt{K^2 + 1} + \ln \left\{ \frac{2}{K^2} + \frac{2}{K^2} \sqrt{K^2 + 1} \right\} \end{aligned}$$

Assume that, for the dividing streamlines, $x \gg y$
when $x \rightarrow \pm \infty$:

So:

$$\frac{\psi_1}{\Gamma/2\pi} = K \frac{y_1}{a} + \ln \left[\left(\frac{y_1}{a} \right)^2 - 1 \right] = \frac{y_1}{a} + \frac{1}{2} \left\{ \ln \left[\left(\frac{x}{a} \right)^2 + \left(\frac{y_1}{a} + 1 \right)^2 \right] + \ln \left[\left(\frac{x}{a} \right)^2 + \left(\frac{y_1}{a} - 1 \right)^2 \right] \right\} \sim$$

$$\sim_{x \rightarrow \pm \infty} K \frac{y}{a} + \frac{1}{2} \left[\ln \left(\frac{x}{a} \right)^2 + \ln \left(\frac{x}{a} \right)^2 \right] =$$

$$= K \frac{y}{a} + 2 \cdot \ln \left(\frac{x}{a} \right) \Rightarrow$$

$$\Rightarrow K \left(\frac{y}{a} - \frac{y_1}{a} \right) = - \ln \left[\frac{\left(\frac{x}{a} \right)^2}{\left(\frac{y_1}{a} \right)^2 - 1} \right] \Rightarrow$$

$$\Rightarrow \frac{y}{a} = \frac{y_1}{a} - \frac{1}{K} \ln \left[\frac{\left(\frac{x}{a} \right)^2}{\left(\frac{y_1}{a} \right)^2 - 1} \right] = \left\{ \frac{y_1}{a} + \frac{1}{K} \ln \left[\left(\frac{y_1}{a} \right)^2 - 1 \right] \right\} - \frac{1}{K} \ln \left(\frac{x}{a} \right)^2$$

$$\frac{\psi_2}{\Gamma/2\pi} = K \frac{y_2}{a} + \ln \left[\left(\frac{y_2}{a} \right)^2 - 1 \right] = K \frac{y}{a} + \frac{1}{2} \left\{ \ln \left[\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} + 1 \right)^2 \right] + \ln \left[\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} - 1 \right)^2 \right] \right\} \sim$$

$$\sim_{x \rightarrow \pm \infty} K \frac{y}{a} + \frac{1}{2} \left[\ln \left(\frac{x}{a} \right)^2 + \ln \left(\frac{x}{a} \right)^2 \right] = K \frac{y}{a} + \ln \left(\frac{x}{a} \right)^2 \Rightarrow$$

$$\Rightarrow K \left(\frac{y}{a} - \frac{y_2}{a} \right) = - \ln \left[\frac{\left(\frac{x}{a} \right)^2}{\left(\frac{y_2}{a} \right)^2 - 1} \right] \Rightarrow$$

$$\Rightarrow \frac{y}{a} = \frac{y_2}{a} - \frac{1}{K} \ln \left[\frac{\left(\frac{x}{a} \right)^2}{\left(\frac{y_2}{a} \right)^2 - 1} \right] = \left\{ \frac{y_2}{a} + \frac{1}{K} \ln \left[\left(\frac{y_2}{a} \right)^2 - 1 \right] \right\} - \frac{1}{K} \ln \left(\frac{x}{a} \right)^2$$

So the dividing streamlines are given by:

$$\frac{y}{a} = \frac{y_1}{a} - \frac{1}{2K} \left\{ \ln \left[\frac{\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} + 1 \right)^2}{\left(\frac{y_1}{a} \right)^2 - 1} \right] + \ln \left[\frac{\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} - 1 \right)^2}{\left(\frac{y_1}{a} \right)^2 - 1} \right] \right\}$$

$$\frac{y}{a} = \frac{y_2}{a} - \frac{1}{2K} \left\{ \ln \left[\frac{\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} + 1 \right)^2}{\left(\frac{y_2}{a} \right)^2 - 1} \right] + \ln \left[\frac{\left(\frac{x}{a} \right)^2 + \left(\frac{y}{a} - 1 \right)^2}{\left(\frac{y_2}{a} \right)^2 - 1} \right] \right\}$$

and their asymptotes for big x are:

$$\frac{y}{a} \sim \left\{ \frac{y_1}{a} + \frac{1}{K} \ln \left[\left(\frac{y_1}{a} \right)^2 - 1 \right] \right\} - \frac{1}{K} \ln \left(\frac{x}{a} \right)^2 = y_{c1} - \frac{1}{K} \ln \left(\frac{x}{a} \right)^2$$

$$\frac{y}{a} \sim \left\{ \frac{y_2}{a} + \frac{1}{K} \ln \left[\left(\frac{y_2}{a} \right)^2 - 1 \right] \right\} - \frac{1}{K} \ln \left(\frac{x}{a} \right)^2 = y_{c2} - \frac{1}{K} \ln \left(\frac{x}{a} \right)^2$$

There are several cases depending on K .

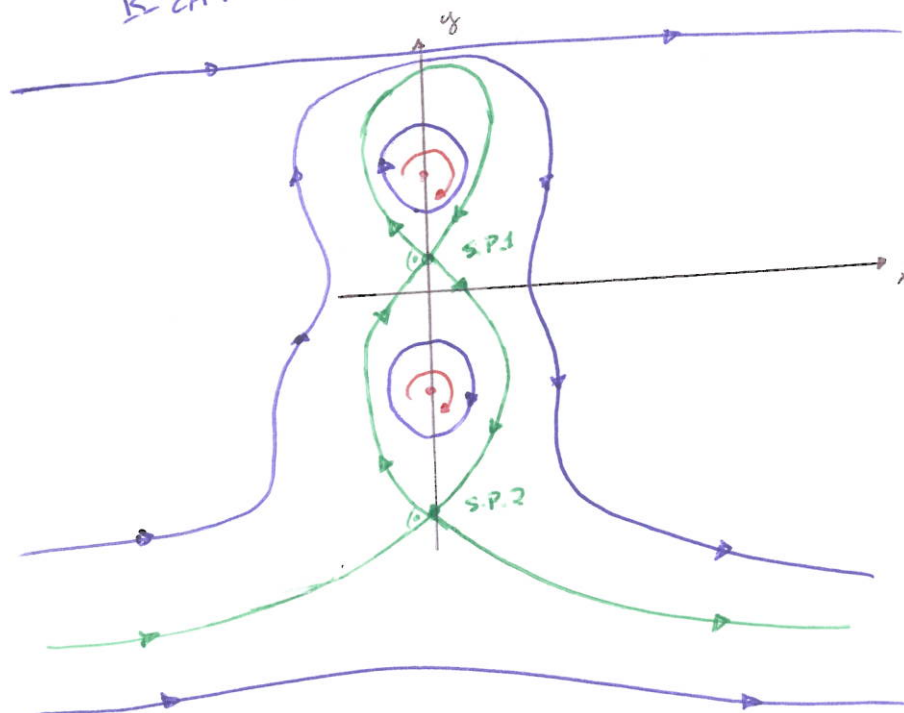
⇒ First, there is a critical case where both dividing streamlines coincide:

$$\frac{\psi_1}{L/2\pi} = \frac{\psi_2}{L/2\pi} \Rightarrow K \left(\frac{y_1}{a} - \frac{y_2}{a} \right) = -\ln \left[\frac{\left(\frac{y_1}{a} \right)^2 - 1}{\left(\frac{y_2}{a} \right)^2 - 1} \right] \Rightarrow$$

$$\Rightarrow \ln \left\{ \frac{\frac{1}{K} - \sqrt{1 + \frac{1}{K^2}}}{\frac{1}{K} + \sqrt{1 + \frac{1}{K^2}}} \right\} + 2K \sqrt{1 + \frac{1}{K^2}} = 0$$

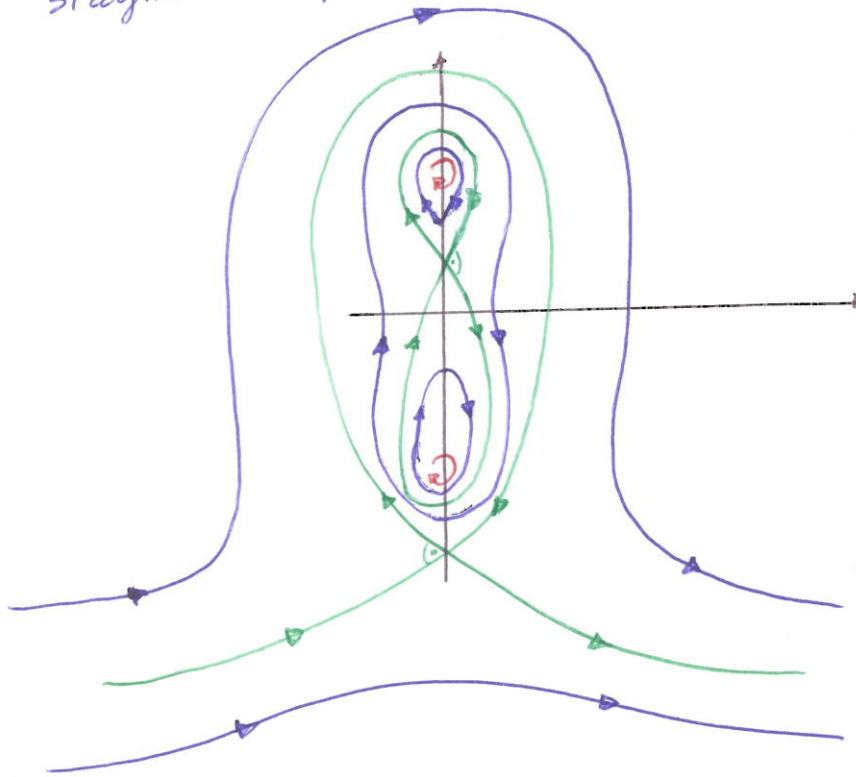
Solving numerically this equations I get:

$$K_{crit} \approx 0.6628$$



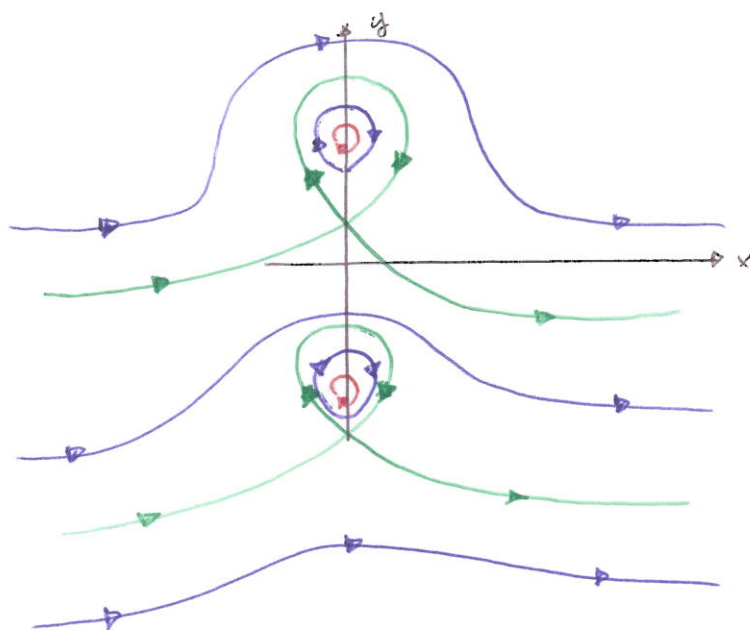
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 ⇒ Second, $R < R_{crit}$, i.e., the vortices are "stronger" than the horizontal stream:

can calculate numerically that the stream line passing through the stagnation point 1 cuts three times the y-axis, and the one passing through the stagnation point 2, cuts only 2:



⇒ Third, $R > R_{crit}$, i.e., the vortices are "weaker" than the horizontal stream.

can calculate numerically that the stream line passing through the stagnation point 1 cuts two times the y-axis, and the one passing through the stagnation point 2, also 2 times:



⇒ There is a subcase in this one. When $R > R_{crit}$ and $y_{c1} > y_{c2}$, we have more solutions. This means that the streamlines pass more than twice through the y-axis.